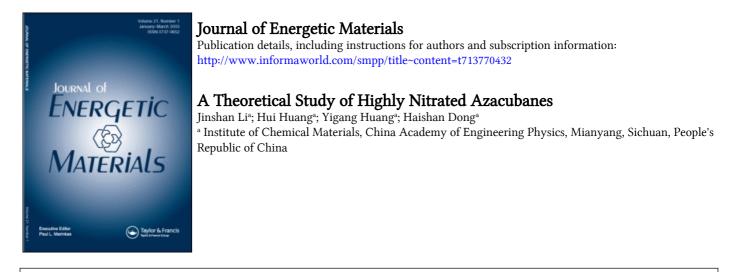
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# A Theoretical Study of Highly Nitrated Azacubanes

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The geometries of highly nitrated azacubanes have been fully optimized at the  $B3LYP/6-31 + G^{**}$  theory level, and the gas phase enthalpies of formation have been obtained by designing the isodesmic reaction in which the azacubane cage skeleton is not destroyed. An efficient procedure for judging the driving force of detonation products of energetic compounds has been developed from the conservation of energy condition. The  $MP2/6-31 + G^{**}//$  $B3LYP/6-31 + G^{**}$  atom-atom overlap-weighted NAO bond order shows that after the full nitration of 1,4-diazacubane and 1,3,5,7-tetraazacubane the weakest C-N bond on the 1,3,5,7-tetraazacubane cage skeleton strengthens slightly, but on the 1,4-diazacubane cage skeleton it weakens. For highly nitrated azacubanes, the introduction of  $-NH_2$  group results in the destabilization of the neighboring C-N bond on the cage skeleton. The shock stability of 2,4,6,8-tetranitro-1,3,5,7-tetraazacubane (TNTAzC) is superior to that of 2,3,5,6,7,8-hexanitro-1,4-diazacubane (HNDAzC). The detonation velocity and pressure for TNTAzC are predicted to reach 11.10 km/s and 101.7 GPa, respectively. The driving force of detonation products of

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HNDAzC is close to that of octanitrocubane, and the driving force of TNTAzC is about 4.0 times as large as that of the widely used high explosive HMX, showing that TNTAzC is a potential candidate of super high-energy compound.

Keywords: 2,3,5,6,7,8-hexanitro-1,4-diazacubane, 2,4,6,8tetranitro-1,3,5,7-tetraazacubane, bond order, density functional B3LYP method, detonation performance

## Introduction

Energetic compounds are widely used [1,2]. There is an ongoing need for high-energy compounds. Although examined as an extremely high-energy compound [3,4], N<sub>8</sub> cubane (octaazacubane) is difficult to synthesize because of its instability. However, if some carbon atoms take the place of nitrogen atoms of octaazacubane (equivalently some carbon atoms of cubane are replaced with nitrogen atoms), the new azacubane molecules will be much more stable than octaazacubane, with the C-N bond being stronger than the N-N bond [5]. On the other hand, because the nitrogen atom in azacubane carries significantly more energy than the carbon atom in cubane, the fully nitrated azacubanes will have better detonation performance with respect to octanitrocubane (ONC), which recently has been synthesized by Eaton and coworkers and predicted to have larger detonation velocity and pressure than the polycyclic nitramine CL-20 [6].

For energetic compounds the detonation performance is a fundamental and important property. In the past, the detonation velocity and pressure (which can be calculated with the Kamlet-Jacobs empirical method [7], the BKW equation of state [EOS] semiempirical method [8], the VLW EOS semiempirical method [9], and so forth) have been used to evaluate theoretically the detonation performance of energetic compounds. But the driving force of detonation products of energetic compounds is sometimes concentrated on, which involves the metal acceleration and can be determined experimentally with the cylinder test and the plane flyer test [2]. The present methods, which use detonation velocity, detonation pressure, or detonation heat to evaluate the driving force of detonation products of energetic compounds, are usually inefficient. Therefore, it is necessary to develop a more efficient theoretical procedure for judging the driving force of detonation products of energetic compounds.

In this article, firstly a procedure for judging the driving force of detonation products of energetic compounds has been developed from the conservation of energy condition. The properties of enthalpy of formation, crystal density, and detonation performance of highly nitrated azacubanes have been investigated from theoretical calculations in order to search for azacubane-based high-energy compounds. Additionally, this work tries to shed some light on the theoretical design of energetic compounds at the ab initio level.

## **Computational Methods and Details**

## Molecular and Crystal Structures

The density functional B3LYP method [10] and the 6- $31 + G^{**}$  basis set are applied to fully optimize the molecular geometries of 1,4-diazacubane (DAzC), 2,3,5,6,7,8-hexanitro-1,4-diazacubane (HNDAzC), 1,3,5,7-tetraazacubane (TAzC), and 2,4,6,8-tetranitro-1,3,5,7-tetraazacubane (TNTAzC) with the Berny gradient approach [11]. Harmonic vibrational frequency has been performed at the  $B3LYP/6-31 + G^{**}$  level, and natural bond orbital analysis [12] performed at the  $MP2/6-31 + G^{**}/B3LYP/6-31 + G^{**}$  level. Considering that the introduction of  $-NH_2$  group into 1,3,5-trinitrobenzene can get insensitive compounds 1,3-diamino-2,4,6-trinitrobenzene (DATB) and 1,3,5-triamino-2,4,6-trinitrobenzene (TATB), 2,7-diamino-3,5,6,8-tetranitro-1,4-diazacubane compounds (DATNDAzC) and 2-amino-4,6,8-trinitro-1,3,5,7-tetraazacubane (ATNTAzC) have also been investigated to discuss the influence of -NH<sub>2</sub> group on the stability of highly nitrated azacubanes. All these calculations have been fulfilled with the GAUSSIAN-03W program [13] using the default Gaussian convergence criteria. The ab initio crystal structure prediction method proposed by Karfunkel and Gdanitz [14] (hereafter called the Karfunkel–Gdanitz method) contained in the Materials Studio software [15] has been employed to predict the crystal density from molecular structure, which proves to be feasible for the density prediction of polynitropyridines [16].

## **Detonation Performance**

For a detonation wave propagating with velocity D into an energetic material initially at rest, the conservation of energy condition is [17]

$$E_1 - E_0 = \frac{1}{2}(P_1 + P_0)(V_0 - V_1)$$
(1)

where  $E_0$ ,  $P_0$ , and  $V_0$  are, respectively, the initial internal energy per unit mass, pressure, and specific volume, and  $E_1$ ,  $P_1$ , and  $V_1$  are, respectively, the final internal energy per unit mass, pressure, and specific volume. Since  $P_1 \gg P_1$ , for explosives Eq. (1) can be written as [17]:

$$E_1 - E_{00} = \frac{1}{2}P_1(V_0 - V_1) + Q \tag{2}$$

where  $E_{00}$  is the specific internal energy of the solid explosive and Q is the specific chemical energy released by the explosive detonation. If the detonation products are regarded as polytropic gases  $p = AV^{-\gamma}$ , where  $\gamma$  is the polytropic exponent, the detonation pressure ( $P_{\rm CJ}$ ) and the specific volume ( $V_{\rm CJ}$ ) at the Chapman-Jouguet (CJ) point can be expressed as [17]:

$$V_{\rm CJ} = \frac{\gamma}{\gamma + 1} V_0 \tag{3}$$

$$P_{\rm CJ} = \frac{1}{\gamma + 1} \rho_0 D^2 \tag{4}$$

$$D = \sqrt{2(\gamma^2 - 1)Q} \tag{5}$$

where  $\rho_0$  is the initial density. The substitution of Eqs. (3), (4), and (5) into Eq. (2) yields the change of internal energy per unit

volume at the CJ point:

$$(E_{CJ} - E_{00})\rho_0 = \frac{\gamma}{\gamma - 1} \frac{D^2 \rho_0}{(\gamma + 1)^2} = \frac{(\rho_0 D^2 - P_{CJ})P_{CJ}^2}{(\rho_0 D^2 - 2P_{CJ})\rho_0 D^2}$$
(6)

Obviously the value of  $(E_{CJ} - E_{00})\rho_0$  can reflect the driving force of detonation products of energetic compounds. Based on the crystal density and the solid-phase enthalpy of formation, the VLW EOS method has been employed to calculate D and  $P_{CJ}$ . The detonation velocity has also been calculated using the convenient Stine method [18]:

$$D = 3.69 + \rho_0 [13.85n_N + 37.74n_N + 68.11n_O + 3.95n_H + 0.1653\Delta_f H_0(s)]/M$$
(7)

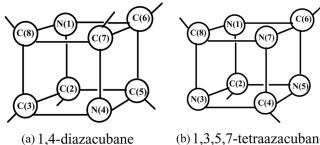
where D is in km/s,  $\rho_0$  is in g/cm<sup>3</sup>,  $\Delta_f H^0(s)$  is the solid-phase enthalpy of formation in kJ/mol, and  $n_{\rm C}$ ,  $n_{\rm N}$ ,  $n_{\rm O}$ , and  $n_{\rm H}$  are, respectively, the numbers of atoms C, N, O, and H, and M is the molecular mass of compound. For CHNO energetic compounds the Stine method is relatively more accurate than the Kamlet-Jacobs empirical method [19].

### **Results and Discussion**

#### Geometry and Bond Order

Molecular structure has an important influence on susceptibility to shock- and impact-induced reactivity [20]. In fact, for most nitroaromatic compounds it has been confirmed that the  $C-NO_2$  bond is of key importance in determining impact and shock sensitivities [21]. In general, the stronger the  $C-NO_2$ bond, the less sensitive the compound. For highly nitrated azacubanes it is realistically postulated that the C-N bonds on the cage skeleton would be related to impact and shock sensitivities due to the existence of large cage strain energy.

The atomic numberings for the cage skeletons of highly nitrated azacubanes are displayed in Fig. 1, which are consistent with the positions of substituents. Vibrational analysis shows that there are no imaginary frequencies for all optimized



(b) 1,3,5,7-tetraazacubane

Atomic numberings for the cage skeletons of highly Figure 1. nitrated azacubanes.

structures, which suggests that all the structures obtained are local minima on the potential energy surface. The calculated longest bond length and the least atom-atom overlap-weighted NAO bond order for the C–C and C–N bonds on the cage skeletons are listed in Table 1. It can be seen from Table 1 that after the full nitration of 1,3,5,7-tetraazacubane the longest C-N bond length on the cage skeleton shrinks by 0.0006 nmand the least C–N bond order increases by 0.006. Therefore, the C–N bond on the cage skeleton of TNTAzC is slightly stronger than that of 1,3,5,7-tetraazacubane. But the full nitration of 1,4-diazacubane makes the C–N bond on the cage skeleton weaken and the C-C bond strengthen. It is worth comparing the bond orders of HNDAC and TNTAzC. For TNTAzC the least C-N bond order is 0.746, which is larger than the least C-N bond order (0.704) of HNDAzC, implying that the shock stability of TNTAzC is superior to that of HNDAzC.

For aromatics and other unsaturated systems, the proper introduction of  $-NH_2$  group can usually get a stable and insensitive molecule. However, for DATNDAzC and ATNTAzC the introduction of  $-NH_2$  group makes the neighboring C-N bond on the cage skeleton elongate significantly. For DATNDAzC,  $r_{2-5}$  reaches 0.1592 nm, and for ATNTAzC,  $r_{1-2}$  reaches 0.1591 nm. Consequently, for highly nitrated azacubanes, the introduction of  $-NH_2$  group has a large destabilizing effect on

The longest bc C-C and C-	ond leng -N bond	th (r/nm) a ls on the ca	The longest bond length $(r/\text{nm})$ and the least atom-atom overlap-weighted NAO bond order $(\beta)$ for the C-C and C-N bonds on the cage skeletons for DAzC, TAzC, and their highly nitrated derivatives	–atom ov DAzC, T.	erlap–weig AzC, and <sup>-</sup>	ghted NAO bond their highly nitra	order (/ ated der	() for the vatives
Compounds	$\mathrm{Par}_{\delta}$	ameters	Compounds Parameters Compounds Parameters Compounds Parameters	Para	meters	Compounds	$\operatorname{Para}$	meters
DAzC	$r_{1-2}$ L	0.1540	DATNDAzC	$r_{4-7}$	0.1574	TNTAzC	$r_{3-8}$	0.1506
	$U_{1-2}$ $\Gamma_{2-3}$	0.1547 0.1547		$p_{4-7}$ $r_{2-5}$	$0.009 \\ 0.1592$	ATNTAzC	$P_{ m 4-5}$ $r_{ m 1-2}$	0.1591
	$eta_{2^{-3}}$	0.805		$eta_{2^{-5}}$	0.790		$eta_{1-2}$	0.665
HNDAzC	$r_{1-8}$	0.1520	TAzC	$r_{1-2}$	0.1512			
	$eta_{1-8}$	0.704		$eta_{1-2}$	0.740			
	$r_{3-8}$	0.1548						
	$\beta_{3-8}$	0.829						

4  $(R) f_{\alpha}$ ÷ 5 JUVD 1-4-1-. ÷ Table 1 . -4+ 2 / -/ 4 Ę کے + -Ē

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the cage skeleton, which is in agreement with a previous ab initio study of nitrocubanes [22]. In summary, for highly nitrated azacubanes it is inappropriate to introduce the  $-NH_2$  group onto the cage skeletons.

## Gas Phase Enthalpy of Formation

For energetic compounds the gas-phase enthalpy of formation  $[\Delta Hf^{0}(g)]$  is an important thermodynamic property. The calculated gas-phase enthalpies of formation for reference compounds CH<sub>4</sub>, CH<sub>3</sub>NH<sub>2</sub>, CH<sub>3</sub>NO<sub>2</sub>, cubane, DAzC, and TAzC using the G3(MP2) method [23] and the atomization scheme are listed in the top lines of Table 2. Table 2 shows that

#### Table 2

Enthalpy and gas-phase enthalpy of formation for CH<sub>4</sub>, CH<sub>3</sub>NH<sub>2</sub>, CH<sub>3</sub>NO<sub>2</sub>, cubane, DAzC, and TAzC at the G3(MP2) level and for HNDAzC, DAHNDAzC, TNTAzC, and ATNTAzC at the B3LYP/6-31++G<sup>\*\*</sup> level at  $P = 1.013 \times 10^5$  Pa and T = 298 K (1 hartree = 2625.50 kJ/mol)

	(	/ /	
Compounds	(kJ/mol)	(kJ/mol)	(kJ/mol)
$CH_4$	$-106,\!118.20$	-74.3	$-74.8^{a}$
$CH_3NH_2$	$-259,\!088.99$	-16.8	$-23.0^{a}$
$CH_3NO_2$	$-642,\!460.40$	-66.1	$-74.7^{a}$
cubane	$-811,\!032.40$	613.9	$622.2^b$
DAzC	$-895,\!219.33$	760.2	
TAzC	$-979,\!432.77$	880.0	
ONC	$-5,\!107,\!233.24$	857.4	
HNDAzC	$-4,\!117,\!703.41$	929.8	
DATNDAzC	$-3,\!334,\!670.17$	769.6	
TNTAzC	$-3,\!128,\!177.20$	1006.6	
ATNTAzC	$-2,\!736,\!676.87$	910.2	

<sup>a</sup>From Weast et al. [29].

<sup>b</sup>From Kybett et al. [30].

the G3(MP2) calculated  $\Delta H_{\rm f}^0(g)$  approaches relatively the corresponding experimental value. The gas-phase enthalpy of formation of azacubane is larger than that of cubane, implying that the replacement of carbon atoms of cubane with nitrogen atoms can boost the energy significantly.

For highly nitrated azacubanes, it is very expensive to use the G3(MP2) method to calculate their enthalpies of formation, so the isodesmic concept has been adopted. The isodesmic reactions are designed as follows, in which the cage skeleton is not destroyed:

$$C_n N_m H_{n-x-y} (NO_2)_x (NH_2)_y + (x+y)CH_4$$
  

$$\rightarrow C_n N_m H_n + xCH_3 NH_2$$
(8)

where  $C_n N_m$  is the cage skeleton and x and y stand for the numbers of  $-NO_2$  and  $-NH_2$  groups, respectively. By employing the B3LYP/6-31 ++ G<sup>\*\*</sup> calculated enthalpies ( $H^0$ ) of all reactants and products, which are corrected for the zero-point energy, the isodesmic reaction, the above G3(MP2) calculated  $\Delta H_f^0(g)$ , and the calculated values of  $\Delta H_f^0(g)$  for highly nitrated azacubanes at  $P = 1.013 \times 10^5$  Pa and T = 298 K are listed in the bottom lines of Table 2. For brevity, the B3LYP/  $6-31 ++ G^{**}$  calculated enthalpies for reference compounds CH<sub>4</sub>, CH<sub>3</sub>NH<sub>2</sub>, CH<sub>3</sub>NO<sub>2</sub>, cubane, DAzC, and TAzC are not displayed in Table 2. It is obvious that highly nitrated azacubanes hold the large enthalpy of formation. Owing to parent contribution, the enthalpy of formation increases greatly from the fully nitrated cubane and diazacubane to tetraazacubane.

#### Crystal Density

Equations (4), (6), and (7) reveal that the crystal density influences directly the detonation performance. During the ab initio prediction of crystal structure, the Dreiding force field [24], which models carbon, nitrogen, oxygen, and hydrogen accurately, has been adopted, and the 17 most common space groups are considered [25]; the asymmetric unit contains one rigid unit. The calculated crystal density ( $\rho_{\rm K-D}$ ) and binding energy ( $\Delta E$ ) by the Karfunkel–Gdanitz method are arranged in Table 3 for the five space groups with lower total energy. For convenient comparison, Table 3 also gives the calculated crystal density ( $\rho_{add}$ ) by the group additivity method [26].

An energetic compound usually has several polymorphs, which can be obtained from recrystallization with different solvents, and we are interested in the polymorph with the largest density. Table 3 shows that for an energetic compound, generally the larger the crystal binding energy, the larger the crystal density. The largest value of  $\rho_{\rm K-D}$  for ONC is 2.065 g/cm<sup>3</sup> with the PI space group, which approaches more the experimental (1.979 g/cm<sup>3</sup>) with respect to the  $\rho_{\rm add}$  (2.111 g/cm<sup>3</sup>). The crystal density for HNDAzC and TNTAzC are predicted to reach 1.975 and 2.156 g/cm<sup>3</sup>, respectively, implying that the two compounds have very large crystal density.

 Table 3

 Calculated crystal density and binding energy for HNDAzC,

 TNAzC, and ONC

		,			
Compounds	$ ho_{ m exp} \ ({ m g/cm}^3)$	$ ho_{ m add} \ ({ m g/cm}^3)$	Space groups	$\Delta { m E} \ ({ m kJ/mol})$	$ ho_{ m K-D} \ ({ m g/cm}^3)$
HNDAzC		2.144	$\begin{array}{c} \mathrm{P2_1} \\ \mathrm{P2_1/c} \\ \mathrm{P1} \\ \mathrm{P1} \\ \mathrm{C2/c} \end{array}$	$142.3 \\137.7 \\133.5 \\133.5 \\132.2$	$     1.975 \\     1.912 \\     1.896 \\     1.897 \\     1.894 $
TNAzC		2.199	$\begin{array}{c} C2/c \\ P2_12_12_1 \\ P2_1 \\ Cc \\ P2_1/c \\ Pca2_1 \end{array}$	$132.2 \\ 149.0 \\ 149.0 \\ 143.5 \\ 138.5 \\ 135.6$	$   \begin{array}{r}     1.894 \\     2.156 \\     2.155 \\     1.996 \\     2.074 \\     2.066   \end{array} $
ONC	1.979	2.111	$\begin{array}{c} P\overline{1} \\ P\overline{1} \\ P2_1 \\ P2_1/c \\ Pna2_1 \\ C2/c \end{array}$	$135.0 \\ 147.3 \\ 134.7 \\ 131.0 \\ 129.7 \\ 128.0$	$2.000 \\ 2.065 \\ 1.954 \\ 1.971 \\ 1.932 \\ 1.930$

## **Detonation Performance**

In order to evaluate the driving force of detonation products of highly nitrated azacubanes, the widely used highly explosive HMX is recommended as a reference compound, and from Eq. (6) the relative driving force of detonation products of other energetic compounds is judged by

$$I_{e} = \frac{(\rho_{0}D^{2} - P_{CJ})P_{CJ}^{2}}{(\rho_{0}D^{2} - 2P_{CJ})\rho_{0}D^{2}} \bigg/ \frac{(\rho_{HMX}D_{HMX}^{2} - P_{HMX})P_{HMX}^{2}}{(\rho_{HMX}D_{HMX}^{2} - 2P_{HMX})\rho_{HMX}D_{HMX}^{2}}$$
(9)

where  $P_{\rm HMX}$  is the experimental detonation pressure (39.0 GPa) for HMX,  $D_{\rm HMX}$  is the detonation velocity (9.11 km/s), and  $\rho_{\rm HMX}$  is the corresponding density (1.894 g/cm<sup>3</sup>) [2]. The experimental  $I_{\rm e}$  by the cylinder test and the calculated  $I_{\rm e}$  by Eq. (9) are listed in Table 4, in which  $E_{\rm cyl}$  is the specific wall kinetic energy at 19 mm wall displacement in the cylinder test and can determine the driving force of detonation products of explosives [2]. Table 4 shows that the calculated  $I_{\rm e}$  by Eq. (9) is very in agreement with the experimental  $I_{\rm e}$ .

In order to calculate the detonation performance of energetic compounds, the solid-phase enthalpy of formation is

 Table 4

 Experimental detonation performance [2] and relative driving force of detonation products for ordinary explosives

Compounds	$ ho_0 \ ({ m g/cm}^3)$	$(\rm km/s)$	$P_{\rm CJ}$ (GPa)	$rac{E_{ m cyl}}{ m (km^2/s^2)}$	$I_e^{ m exp}$	$I_e^{ m cal}$
$CH_3NO_2$	1.130	6.35	12.5	0.745	0.43	0.35
TNT	1.63	6.93	21.0	0.975	0.56	0.58
TATB	1.86	7.66	25.9	1.045	0.62	0.64
PETN	1.765	8.26	33.5	1.575	0.90	0.96
RDX	1.767	8.70	33.8	1.600	0.92	0.88
BTF	1.859	8.49	36.0	1.680	0.96	0.99
HMX	1.894	9.11	39.0	1.745	1.00	1.00

necessary. A procedure for converting gas to solid-phase enthalpy of formation has been developed by Politzer et al. [27] on the basis of the calculated electrostatic potential on the molecular surface. In this work,  $\Delta H_f^0(s) = \Delta H_f^0(g) - \Delta E$ . Based on  $\rho_{\text{K-D}}$  and  $\Delta H_f^0(s)$ , the calculated values of D,  $P_{\text{CJ}}$ , and  $I_e$  for HNDAzC, TNAzC, and ONC are summarized in Table 5. For ONC and HNDAzC, the calculated D by the Stine method is 0.15–0.35 km/s larger than by the VLW EOS method, but for compound TNTAzC is 0.32–0.90 km/s less than by the VLW EOS method. As a whole, the detonation velocities from the two methods are in good agreement. The detonation velocity and pressure of HNDAzC are close to the

Table 5Calculated solid phase enthalpy of formation, detonationvelocity, detonation pressure, and relative driving force ofdetonation products for HNDAzC, TNAzC, and ONC

				V	ΊLW	
Compounds	Space groups	$\Delta H_f^0({ m s}) \ ({ m kJ/mol})$	$\begin{array}{c} \text{Stine} \\ D(\text{km/s}) \end{array}$	$D(\rm km/s)$	$P_{\rm CJ}$ (GPa)	$I_{ m e}$
HNDAzC	$\begin{array}{c} P2_1 \\ P2_1/c \\ P1 \end{array}$	787.5 792.1 796.3	9.82 9.62 9.58	9.47 9.37 9.34	58.2 53.8 52.9	1.98 1.78 1.75
TNTAzC	$\begin{array}{c} P1 \\ C2/c \\ P2_{1}2_{1}2_{1} \\ P2_{1} \end{array}$	796.3 797.6 857.6 857.6	$9.58 \\ 9.57 \\ 10.68 \\ 10.67$	$9.34 \\ 9.34 \\ 11.10 \\ 11.10$	52.9 52.7 101.7 101.7	$     1.75 \\     1.73 \\     4.02 \\     4.02 $
ONC	$\begin{array}{c} \mathrm{Cc} \\ \mathrm{P2}_{1}/\mathrm{c} \\ \mathrm{Pca2}_{1} \\ \mathrm{P\overline{1}} \end{array}$	863.1 868.1 871.0 710.1	$10.16 \\ 10.42 \\ 10.40 \\ 9.91$	$     \begin{array}{r}       11.06 \\       11.08 \\       11.08 \\       9.61     \end{array} $	96.7 99.2 100.1 60.7	3.96 3.99 4.08 2.00
0110	$\begin{array}{c} P2_1 \\ P2_1/c \\ Pna2_1 \\ C2/c \end{array}$	710.1 $722.7$ $726.4$ $727.7$ $729.4$	$9.59 \\ 9.64 \\ 9.52 \\ 9.52$	$9.41 \\ 9.44 \\ 9.37 \\ 9.36$	53.6 54.7 52.4 52.3	1.72 1.76 1.67 1.67

corresponding values of ONC. From the three compounds, TNTAzC has the largest detonation velocity and pressure, which are 11.10 km/s and 101.7 GPa, respectively.

The value of  $I_{\rm e}$  shows that HNDAzC and ONC both have the close driving force of detonation products, which is 1.7– 2.0 times as large as that of HMX, implying that HNDAzC is more powerful than the most powerful used nonnuclear explosive CL-20, whose driving force of detonation products exceeds that of HMX by about 14% [28]. The driving force of detonation products of TNTAzC is about 4.0 times as large as that of HMX, showing that TNTAzC is a potential candidate of super high-energy compound.

#### Conclusions

This work has investigated the geometries, electronic structures, enthalpies of formation, crystal densities, and detonation parameters of highly nitrated azacubanes. An efficient theoretical procedure for judging the driving force of detonation products of energetic materials has been developed from the conservation of energy condition. After the full nitration of 1.4-diazacubane and 1,3,5,7-tetraazacubane, the weakest C-N bond on the 1,3,5,7-tetraazacubane cage skeleton strengthens slightly but on the 1.4-diazacubane cage skeleton weakens. For highly nitrated azacubanes, the introduction of  $-NH_2$  group results in the destabilization of the neighboring C-N bond on the cage skeleton. The MP2/6-31++ $G^{**}//$  $B3LYP/6-31 ++ G^{**}$  atom-atom overlap-weighted NAO bond order implies that the shock stability of TNTAzC is superior to that of HNDAzC. The values of crystal density for HNDAzC and TNTAzC can reach 1.975 and  $2.156 \,\mathrm{g/cm^3}$ , respectively. The detonation velocity and pressure for TNTAzC are predicted to be 11.10 km/s and 101.7 GPa, respectively. The driving force of detonation products of HNDAzC is close to that of ONC, and the driving force of TNTAzC is about 4.0 times as large as that of HMX, showing that TNTAzC is a potential candidate of super high-energy compound.

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## References

- Gibbs, T. R. and A. Popolato. 1980. LASL High Explosives Property Data. Berkeley, CA: University of California Press.
- [2] Dobratz, B. M. and P. C. Crawford. 1985. LLNL Explosives Handbook—Properties of Chemical Explosives and Explosive Simulants. California, USA: Lawrence Livermore National Laboratory.
- [3] Leininger, M. L., T. J. van Huis, and H. F. Shaefer, III. 1997. Nitration of 2-substituted pyrimidine-4,6-diones, structure and reactivity of 5,5-gem-dinitropyrimidine-4,6-diones. *Journal of Physical Chemistry A*, 101: 4460.
- [4] Bruney, L. Y. and D. L. Strout. 2003. Are four oxygens enough to stabilize the N(sub)8 cubane? *Journal of Physical Chemistry A*, 107: 5840.
- [5] Politzer, P., P. Lane, and J. J. M. Wiener. 1999. Carbocyclic and Heterocyclic Cage Compounds and Their Building Blocks. Stanford, CT: JAI Press.
- [6] Zhang, M. X., P. E. Eaton, and R. Gilardi. 2000. Hepta- and octanitrocubanes. Angewandte Chemie International Edition, 39: 401.
- [7] Kamlet, M. J. and S. J. Jacobs. 1968. Chemistry of detonation I: A simplified method for calculation detonation properties of CHNO explosives. *Journal of Chemical Physics*, 48: 23.
- [8] Mader, C. L. 1979. Numerical Modeling of Detonation. Berkeley, CA: University of California Press.
- [9] Wu, X. 1985. A study on thermodynamic functions of detonation products. Propellants, Explosives, Pyrotechnics, 10: 47.
- [10] Becke, A. D. 1993. Density-functional thermochemistry III: The role of exact exchange-correlation functions. *Journal of Chemical Physics*, 98: 5648.
- [11] Schlegel, H. B. 1982. Optimization of equilibrium geometries and transition structures. Journal of Computational Chemistry, 3: 214.

- [12] Reed, A. E., B. Robert, and F. Weinhold. 1985. Natural population analysis. *Journal of Chemical Physics*, 83: 735.
- [13] Frisch, M. J., G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, V. G. Zakrzewski, J. A. Pople, et al. 2003. *Gaussian 03.* Rev. B.04. Pittsburgh, PA: Gaussian, Inc.
- [14] Karfunkel, H. R. and R. J. Gdanitz. 1992. Ab initio prediction of possible crystal structures for general organic molecules. *Journal* of Computational Chemistry, 13: 1171.
- [15] MS Modeling. Ver. 4.00. 2005. San Diego: Accelrys Software Inc.
- [16] Li, J. S., Y. G. Huang, and H. S. Dong. 2005. A theoretical study of polynitropyridines. *Journal of Energetic Materials*, 23: 133.
- [17] Fickett, W. and W. C. Davis. 2000. Detonation Theory and Experiment. New York: Dover Publications.
- [18] Stine, J. R. 1990. On predicting properties of explosives detonation velocity. *Journal of Energetic Materials*, 8: 41.
- [19] Li, J. S. 1994. A novel method for calculating detonation velocity. *Chinese Journal of Explosives and Propellants*, 17(3): 28 (in Chinese).
- [20] Politzer, P. and S. Boyd. 2002. Molecular dynamics simulations of energetic solids. *Structural Chemistry*, 13: 105.
- [21] Owens, F. J., K. Jayasuriya, L. Abrahmsen, and P. Politzer. 1985. Computational analysis of some properties associated with the nitro groups in polynitroaromatic molecules. *Chemical Physics Letters*, 116: 434.
- [22] Murray, J. S., J. M. Seminario, and P. Politzer. 1991. Effect of the amine group on relative bond strengths in cubane and azacubanes. *Structural Chemistry*, 2: 153.
- [23] Curtiss, L. A., P. C. Redfern, K. Raghavachari, V. Rassolov, and J. A. Pople. 1999. Gaussian-3 theory using reduced Møller-Plesset order. *Journal of Chemical Physics*, 110: 4703.
- [24] Mayo, S. L., B. D. Olafson, and W. A. Goddard. 1990. DREID-ING: A generic force field for molecular simulations. *Journal of Physical Chemistry*, 94: 8897.
- [25] Belsky, V. K. and P. M. Zorkii. 1977. Distribution of organic homomolecular crystals by chiral types and structural classes. *Acta Crystallographica A*, 33: 1004.
- [26] Stine, J. R. 1981. Prediction of Crystal Densities by Group Additivity. Albuquerque, NM: Los Alamos National Laboratory.
- [27] Politzer, P., Y. Ma, P. Lane, and M. C. Concha. 2005. Computational prediction of standard gas, liquid, and solid-phase heats of

formation and heats of vaporization and sublimation. International Journal of Quantum Chemistry, 105: 341.

- [28] Simpson, R. L., P. A. Urtiew, D. L. Ornellas, G. L. Moody, K. J. Scribner, and D. M. Hoffman. 1997. CL-20 performance exceeds that of HMX and its sensitivity is moderate. *Propellants*, *Explosives*, *Pyrotechnics*, 22: 249.
- [29] Weast, R. C., D. R. Lide, M. J. Astle, and W. H. Beyer. 1990. CRC Handbook of Chemistry and Physics. 70th ed. Boca Raton, FL: CRC Press.
- [30] Kybett, B. D., S. Carroll, P. Natalis, D. W. Bonnell, J. L. Margrave, and J. L. Franklin. 1966. Thermodynamic properties of cubane. *Journal of the American Chemical Society*, 88: 626.